Multiple Intra-Tube Junctions in the Inner Tube of Peapod-Derived Double Walled Carbon Nanotubes: Theoretical Study and Experimental Evidence

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The coalescence process of fullerenes in the hollow core of single wall carbon nanotube is systematically explored by kinetic Monte Carlo method. Two elongation (or growth) modes via the coalescence (i) between an inner tube and fullerenes and (ii) between neighboring inner tubes are identified. It is found that the coalescence of two inner tubes mostly creates a very stable intra-tube junction which is composed of multiple pentagon-heptagon pairs. As a consequence, the study predicts that the inner tube of peapod-grown double walled carbon nanotubes (DWNTs) must contain many intra-tube junctions. Careful high resolution transmission electron microscopy observation on peapod-grown DWNT sample provides an experimental evidence on the presence of the junctions.

1. Introduction

Carbon peapods, the hybrids of two important carbon nanomaterials (i.e., carbon nanotubes (CNTs) and fullerenes), have drawn great attention since their first experimental synthesis in 1998.\textsuperscript{1-7} It is amazing that the coalescence of the encapsulated fullerenes leads to the formation of uniform inner tube and then transforms a peapod into a double-walled CNT (DWNT).\textsuperscript{8,9} After 10-year extensive study, “peapod” method becomes a standard way of synthesizing DWNTs. Compared to other synthetic methods, “peapod” method has several advantages: (i) The process is free of catalyst; (ii) It can provide small-diameter DWNT sample;\textsuperscript{10} (iii) The products have a very high DWNT to SWNT ratio.\textsuperscript{11}

The theoretical study of the coalescence of encapsulated fullerenes has been a great challenge for many years.\textsuperscript{5, 12-24} In terms of the methods used, the theoretical studies previously published can be classified into three categories:

Molecular dynamics (MD) simulation, including both classical MD and quantum MD, is one of the most common methods.\textsuperscript{13, 26-31} It’s very unfortunate that none of previous MD simulations has been able to end up with a well-defined chirality of inner tube and thus the information about the peapods derived DWNTs is incomplete. This is due to the limited time scale of MD simulation. The time scale for classical MD and for quantum MD is ns and ps, respectively whereas the experimental time is ranging from minutes to hours. Their difference is more than 12 orders of magnitude.

Kinetic simulations based on C-C bond rotation or so called the generalized Stone-Wales (GSW) transformation have been widely used to simulate the coalescence process of fullerenes.\textsuperscript{32} The kinetic simulations are very diverse compared to the MD simulation. In 2002, Zhao and coworkers have discovered a manual bond rotation procedure to transform a C\textsubscript{60} dimer into a short C\textsubscript{70} SWNT\textsuperscript{33} and then the method has been applied to study the welding between two SWNTs and between a SWNT and a fullerene.\textsuperscript{34, 35} This manual process is very time consuming and thus can only be used to explore a very limited number of samples. Furthermore this procedure is guided by several geometrical rules, which is very different from the real driving force in a natural process, the change of energy and the energy barrier. Latterly, the transition state barrier of every bond rotation has been calculated by Han et al..\textsuperscript{36}

In 2006, the action-derived molecular dynamics (ADMD) method was introduced to simulate the fullerene coalescence process by Lee and coworkers.\textsuperscript{37} The energy difference is the driving force in a ADMD simulation and thus the ADMD simulation is able to produce a reasonable low energy path for fullerene coalescence. While, same as the manual work, the final structure of the ADMD simulation is predetermined. Thus the information about inner tube structure, e.g., the chiral angle or diameter correlation between the inner and outer tubes, can’t be retrieved.

We recently have developed an energy driven kinetic Monte Carlo (EDKMC) method to study the structural transformation of nano carbon materials.\textsuperscript{38, 41} In a EDKMC trajectory, the energy change of a C-C bond rotation is used as the driving force. The EDKMC method has been successfully applied to simulate the coalescence of two fullerenes in a SWNT recently.\textsuperscript{38} The study shows that there is no chiral angle correlation between the inner and outer tubes and the initial inner tube always tends to have a large chiral angle which is near 30 degree.\textsuperscript{38} The agreement between these theoretical predictions and experimental measurements validated the EDKMC method.\textsuperscript{38} Beyond, previous study has shown that the space constrain of outer tube plays a crucial role in the inner tube formation. The final inner tube
always tends to fit the outer tube with a wall-wall distance range of 0.34-0.38 nm, which is exact the same as the experimentally observed wall-wall distance of DWNTs.42

In this article, we present a systematic EDKMC study of inner tube elongation. Two processes of inner tube elongation, namely inner tube + C_{60} coalescence and inner tube + inner tube coalescence are identified. For the inner tube + C_{60} coalescence, most inner tubes tend to maintain its original chirality and the longer the tube stem is, the robust er its chirality is. For the inner tube + inner tube process, an inner tube with a longer chirality identified stem has larger probability for its chirality to be inherited. An intra-tube junction, which is composed of multiple pentagon-heptagon pairs (5|7)s, will be formed between two longer inner tubes. This study reveals that the inner tube of a peapod-derived DWNT must contain highly concentrated intra-tube junctions in the inner wall. Furthermore, careful high resolution transmission electronic microscopy (HRTEM) observation was carried out in order to support both the tube elongation processes and the formation of the intra-tube junctions.

2. Simulation Method

The EDKMC method is used to study the elongation process of short inner tubes. Denoting the energy of the relaxed initial structure as \( E_i \), the energy of final structure which is obtained by rotating a randomly selected C-C bond of the initial structure by 90 degree as \( E_f \). Similar to the standard Monte Carlo simulation, the probability of accepting the final structure is 1 if \( \Delta E = E_f - E_i < 0 \) or \( \exp(-\Delta E/kT) \) if \( \Delta E > 0 \), where \( k \) is the Boltzmann constant and \( T \) is the temperature. More details of the EDKMC method and its validity are shown in previous publication.38 In previous studies, we have shown that there’s no tube chiral angle correlation between the inner wall tube and outer wall tube and thus a (18, 0) SWNT was chosen as the only outer wall tube in all trajectories. The Monte Carlo temperature of all trajectories was set as 2000 K, which is close to the experimental temperature of deriving peapods into DWNTs.

For the popular C_{60}s@SWNT peapods, the formation of peapods derived DWNT starts with a C_{60} dimer. The fusion of a C_{60} dimer resulting in a short C_{120} SWNT, C_{60} + C_{60} \rightarrow C_{120}
SWNT, has been extensively studied before. Here we consider the inner tube elongation from a short C_{120} SWNT. In the elongation process, if a C_{120} is surrounded by C_{60}s, the next step must be the fusion of the C_{120} with a nearby C_{60}, C_{120} + C_{60} \rightarrow C_{180}, and then, C_{180} + C_{60} \rightarrow C_{240}⋯. Beyond the continuous growth of an inner tube by adsorbing C_{60}s, the meeting of two tubes and their fusion is another important process towards inner tube elongation, e.g., C_{120}+C_{120}, C_{120}+C_{180} or C_{180}+C_{180}.

So, in order to achieve a comprehensive understanding of the elongation process, we use C_{120}s and C_{180}s to represent short tube and long inner tube, respectively and designed five different inner tube elongation processes, which are (i) C_{120} + C_{60}; (ii) C_{180} + C_{60}; (iii) C_{120} + C_{120}; (iv) C_{120} + C_{180} and (v) C_{180} + C_{180}, respectively.

In all simulations, the initial C_{120}s are obtained by the coalescence of two C_{60}s and C_{180}s are formed in the coalescence of C_{120} + C_{60}. Certainly, the fusion of further longer tubes (e.g., C_{240}) can be considered as well. But as we can see, the robusticity of the C_{180} inner tubes ensures that it’s not necessary to consider longer inner tubes beyond C_{180}.

3. Results

3.1 The elongation of short inner tubes: C_{120} + C_{60}

Fig. 1a-f depict the elongation of a short (6,5) SWNT by adsorbing a C_{60}. The EDKMC simulation starts with an initial configuration that the short SWNT is linked to the C_{60} by two C-C covalent bonds, which can be viewed as an extreme narrow neck (Fig. 1a). During the initial stage, most GSW bond rotations occur around the neck area because of fast relaxation of the high curvature energy. The curvature energy drops notably (Fig. 1h) and the neck becomes thicker and thicker (Fig. 1a→b→c→d→e). Eventually, a perfect longer (6,5) tube was formed after about 350 GSW steps (Fig. 1f and Fig. 1g). After that, the energy only fluctuates in a very small amplitude (Fig. 1h) and the stem of the tube remains unchanged. In this trajectory, the chirality identified stem of the short SWNT is maintained and thus the final inner tube inherits the original chirality of the short tube. This simulation presents the robusticity of the short inner tubes.

In order to further test the robusticity of inner tubes, the coalescences of five different C_{120} tube isomers, which are (6,5), (5,5), (10,0), (6,3) and (9,0) tubes, with a C_{60} are systematically studied. 20~40 EDKMC successful trajectories were run for every combination (details to see Table S1). The statistics over the final inner tubes’ chiral angles and wall-wall distances are shown in Fig. 1i and 1j. It can be clearly seen that all tubes maintain their original chiralities in the EDKMC trajectories. For example, more than half of the final tubes obtained from the coalescences of (9,0) inner tube with C_{60} are chiral tubes. In contrast, short inner tubes with large chiral angles tend to maintain their original chiralities, e.g., 83% of the (5,5) tubes remain the same after the fusion with a C_{60}. The robusticity of large chiral angle tubes is in agreement with what is shown in previous study.

3.2 The elongation of long inner tubes: C_{180} + C_{60}

Fig. 2a-f are snapshots taken in the coalescence of a (8,4) C_{180} inner tube and a C_{60}. During the coalescence, several pentagon-heptagon pairs (5|7s) are excited from the left end of the C_{180} tube and these 5|7s gradually propagate through the tube wall towards the inner tube-C_{60} neck area and sink there (b→c→d). The final inner tube, whose chirality is (6,5), is notably smaller than the original tube.
Fig. 3. (a-e) Snapshots in the fusion of two (9,0) C_{120} inner tubes into a (5,5) C_{120} tube. (f) Five different C_{120} tubes observed in different trajectories.

(Fig. 2f). This indicates that the wall-wall distance between the inner and outer tubes is a crucial parameter that controls the robusticity of the inner tube. The wall-wall distance between the (8,4) inner tube and the (18,0) outer tube is only 0.29 nm (Fig. 2h), which is 18% smaller than the equilibrium layer-layer distance in graphite, ~0.34 nm. On the other hand, the final (6,5) tube, which has a wall-wall distance of 0.355 nm, fit the (18,0) tube perfectly.

Except the large diameter (8,4) C_{180} inner tube, other three inner tubes, (9,0), (7,3) and (6,5), explored in this study have well fitted wall-wall distances with the (18,0) outer tube, which are 0.352 nm, 0.331 nm and 0.355 nm, respectively. As a consequence, every trajectory (see Table S2†) results in a tube with exact same chirality as the original one’s (Fig. 2g). The apparent contrast with the results shown in Fig. 1 indicates that an inner tube with a chirality identified stem of 2 nm or longer is very robust and its chirality is very hard to be changed unless its wall-wall distance with the outer tube is far from the equilibrium distance, 0.34 nm.

3.3 The competition between two short tubes: C_{120} + C_{120}

Here we focus on the fusion of two short inner tubes, C_{120} + C_{120}. Fig. 3a–e are representative snapshots during the fusion of two (9,0) C_{120} tubes. At the initial elongation stage, the inner tube on the right side becomes very defective and then it turns into a (5,5) tube (Fig. 3a→b→c). A intra-tube junction, which contains a few 5|7s and connects the (9,0) tube on the left and the newly formed (5,5) tube on the right, is formed during the process (Fig. 3c). In the following simulation, the 5|7 pairs around the intra-tube junction gradually glide to the left side of the inner tube and disappeared into the tube cap eventually (Fig. 3c→d→e). The full EDKMC trajectory costs about 2000 GSW steps, which is about one order of magnitude slower than the fusion of C_{120} + C_{60}. The long trajectory implies that the fusion of two short tubes into a single chirality one is very difficult and the diffusion of a tube-tube junction is very slow.

It’s important to note that the final tubes can be very diverse. Beside the before mentioned (5,5) tube, other chiral tubes, for example (9,0), (8,1), (8,2), (7,3) tubes, have been observed in EDKMC trajectories (see Fig. 3f).

3.4 The competition between a longer inner tube and a shorter inner tube: C_{180} + C_{120}

From the results presented above, we can conclude that these very short C_{120} tubes cannot maintain their chiralities during the fusions of C_{120} + C_{60} or C_{120} + C_{120} but the longer C_{180} tubes are very robust. So it’s very natural to suspect that a longer tube is going to beat a short one to inherit its gene, the chirality, in the fusion of C_{180} + C_{120}. Fig. 4 visibly shows such a trend, the stem of the (9,0) tube becomes longer and longer and eventually dominates the final tube in the coalescence of a (9,0) C_{120} + (6,5) C_{120} tube. This indicates that the tube length is the most critical parameter in the tube-tube fusion.

3.5 The competition between two longer tubes: C_{180} + C_{180}

So far, we can conclude that it’s very hard to change the chirality of an inner tube whose chirality identified stem is longer than 2 nm. Therefore, what will happen if two long C_{180} inner tubes meet each other and fuse into one? The fusion process of a (5,5) C_{180} tube and a (9,0) C_{180} tube is shown in Fig. 5. A tube-tube neck which has nearly same diameter as the original tubes is formed in between after ~500 GSW annealing steps (a→b→c). An intra-molecular ring of defects that is composed of a few 5|7 pairs can be identified around the neck area (c). The position of the ring does not change in the following long EDKMC simulation (up to 3000 GSW steps) though the arrangement of the 5|7s becomes different (c→d). The robusticity of the two inner tubes results in a stable intra-tube junction. To link two different SWNTs into one, such a junction is inevitable. Similar intra-tube junctions are seen in other trajectories. As shown in Fig. 5f, the junctions of (9,1)|(9,0), (7,3)|(9,0) and (6,4)|(9,0) contains 1, 3 and 4 5|7 pairs, respectively.

4. Discussion

4.1 The intra-tube junction as a rolled up graphene grain boundary

As shown in the perspective view of the intra-tube junction between a (5,5) tube and a (9,0) tube (Fig. 5e), the 5|7 pairs form a ring of defects in a regular manner, the pentagon of a 5|7 pair contacts with the heptagon of a neighboring pair or pentagons and heptagons appear alternatively along the tube circumference. The novel ring formation is a consequence of minimizing the junction’s formation energy. Recent study in graphene grain boundary reveals that linear arrangement of 5|7 pairs minimizes the boundary formation energy because of the cancellation of the compression and tension strains around a 5|7 pair. A carbon nanotube can be viewed as a rolled up graphene and therefore the ring of 5|7s can be viewed as a rolled up graphene grain boundary. Fig. 5g shows the number of 5|7 pairs of an intra-tube junction
as a function of chiral angle difference between the two tubes. A
clear linear relationship, $N_{5|7} \sim \Delta \theta$, is presented. Such a behavior
is a bit surprising since it has been theoretically proved that only
one pentagon and one heptagon are sufficient to connect two
arbitrary SWNTs (Fig. 6a $\rightarrow$ b). It’s important to note that the
previous theory allows the hybrid tube wall to turn a sharp angle
at the junction (Fig. 6b). While, a peapod derived inner tube is
constrained inside a straight outer tube and thus has to maintain
its straightness at the junction (Fig. 6d). So the previous ideal
model of tube-tube junction can’t be applied to understand the
formation of peapod derived DWNTs. Here we propose another
model that the intra-tube junction should be viewed as a rolled up
graphene grain boundary to explain the multiple $5|7$ formation. In
that model, a straight graphene nanoribbon (GNR) is cut along a
direction perpendicular to the grain boundary (Fig. 6c) and then
been folded into a tube (Fig. 6c $\rightarrow$ d). The straight GNR edges
ensure the straightness of the folded tube (Fig. 6d).

As has been discussed in detail in recent literatures, a $5|7$ in
graphene is an edge dislocation in a 2D hexagonal lattice and the
magnitude of its Burgers vector is $b = 0.246 \text{ nm}$. In a graphene
grain boundary, the requirement of formation energy minimiza-
tion aligns all the $5|7$ pairs linearly. So the mismatch angle
between both sides can be estimated by the following equation: $^{47}$

$$\sin(\Delta \theta) = b/d$$  \hspace{1cm} (1)

where $\Delta \theta$ is the mismatching angle and $d$ is the average distance
between neighboring $5|7$ pairs. As a rolled up grain boundary, the
number of $5|7$ pairs in an intra-tube junction is

$$N_{5|7} = l/d = l \times \sin(\Delta \theta)/b \sim l \times \Delta \theta/b$$  \hspace{1cm} (2)

![Fig. 5](image-url) (a-d) Some typical snapshots during the coalescence of a (9,0) C$_{180}$ inner tube and a (5,0) C$_{180}$ inner tube. (e) The perspective view of the tube-tube junction. (f) Several other tube-tube junctions formed in the coalescence of a (9,0) C$_{180}$ tube and a few other C$_{180}$ tubes. (g) The number of $5|7$ pairs at the tube-tube junction as a function of chiral angle difference between two tubes.

![Fig. 6](image-url) Two arbitrary tubes can be connected by a pentagon and a
heptagon. As an example, a (9,0) tube and a (5,5) tube is connected by
separated pentagon and heptagon (a-b) but the tube wall turns a sharp angle of ~ 30 degree at the intra-tube junction (a). (c)$\rightarrow$(d) Rolling a graphene grain boundary into a straight tube with an intra-tube junction. The grain boundary contains multiple $5|7$ pairs forming a ring of
topological defects (e).
Fig. 7 Transmission electron microscope images of peapod-grown DWNTs that are grown by thermally treating peapods at 1700°C in argon using a graphite furnace at (a) low- and (b-d) high-resolutions.

where \( l \) is the length of tube circumference, \( \Delta \theta \) here is the chiral angle difference between the two tube sections. For those inner tubes encapsulated inside a (18,0) SWNT, their diameters are around 0.7 nm. The number of 5|7 pairs is \( N_{57} \sim 0.15 \times \Delta \theta \) (in degree). For a junction between (5,5) and (9,0) SWNTs, we have \( \Delta \theta = 30 \) degree and \( N_{57} = 4.5 \), which fits the simulated number \( N_{57} = 5 \) perfectly (Fig. 5g).

4.3 A growth model of peapod-grown DWNT

A model of peapod-derived DWNT is shown in Fig. 8 to summarize the study. At the initial nucleation stage, the simultaneous coalescence of many fullerenes in a long carbon peapods results in many isolated short inner tubes of different chiralities inside (Fig. 8 a→b). These short inner tubes grow longer and longer by exhausting neighboring fullerenes (Fig. 8 b→c). Once fullerenes between two inner tubes exhausted, the two tubes contact each other directly and then fuse into one straight tube that contains a stable intra-tube junction (Fig. 8d). The final long inner tube which is formed

by the fusion of many short ones possesses many intra-tube junctions (Fig. 8 e).

5. Conclusions

In conclusion, two inner tube elongation processes in the formation of peapods derived DWNTs are identified by the EDKMC simulation: (i) fusion of tube + C60 and (ii) coalescence of tube + tube. In both processes, only the chirality of very short inner tubes whose stem length less than 2 nm might be changed and a long inner tube shows its extreme robusticity because of the confinement effect of the outer tube. Our study predicates that the fusion of two long inner tubes must result in a robust intra-tube junction. The formation of the intra-tube junction, which can be viewed as a folded graphene grain boundary, contains multiple pentagon-heptagon pairs (5|7) forming a defect ring. The number of the 5|7s in a junction is proportional to the product of the inner tube’s chiral difference and the length of the inner tube circumference. Beyond, careful HRTEM imaging on peapods derived DWNTs samples clearly represents both coalescence processes and shows solid evidence of the intra-tube junction.

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